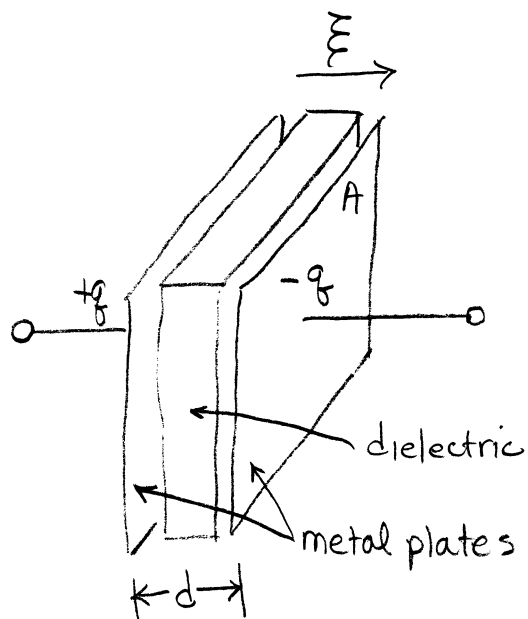


Chapter 6 Capacitance and Inductance



For small d the electric field is given by

$$\vec{E}(t) = \frac{q(t)}{\epsilon A}$$

$$\text{but } \vec{E}(t) = \frac{v_c(t)}{d}$$

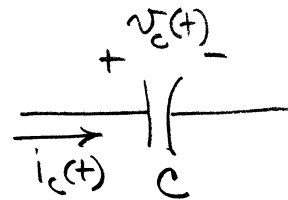
Solving for q
$$q(t) = \epsilon A \vec{E}(t) = \underbrace{\frac{\epsilon A}{d}}_{C} v_c(t)$$

This is a constant dependent upon the physical construction called the capacitance. C

$$q(t) = C v_c(t)$$

Differentiating

$$i_c(t) = \frac{dq(t)}{dt} = C \frac{dv_c}{dt}$$



This the the i - v characteristic of the capacitor.

Consider what this means

- ① when $v_c = \text{constant (dc)}$ the current is zero
- ② discontinuous voltage like $\delta(t)$ would require and infinite current, Because of this voltage across a capacitor must be continuous.

Integrate $i_c(t) = C \frac{dv_c}{dt}$

$$\frac{1}{C} i_c(t) dt = dv_c$$

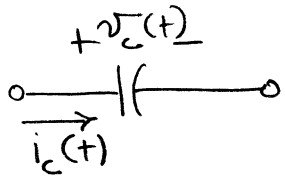
$$\frac{1}{C} \int_{t_0}^t i_c(x) dx = \int_{v_c(t_0)}^{v_c(t)} dv_c = v_c(t) - v_c(t_0)$$

$$\therefore v_c(t) = v_c(t_0) + \frac{1}{C} \int_{t_0}^t i_c(x) dx$$

In most of our problems we will

① define $t_0 = 0$

② recognize $v_c(t_0)$ as the initial voltage on the capacitor.



Consider the power $p_c(t) = i_c(t) v_c(t)$

$$= \left[C \frac{dv_c}{dt} \right] v_c(t)$$

$$p_c(t) = \frac{d}{dt} \left[\frac{1}{2} C v_c^2(t) \right]$$

This is VERY interesting.

$p_c(t) > 0$ capacitor absorbing power

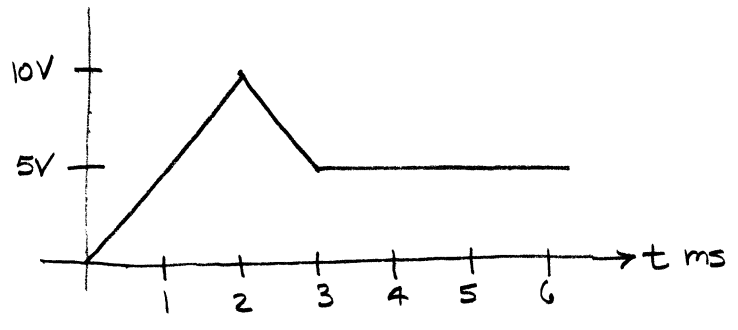
$p_c(t) < 0$ capacitor releasing stored power

Since $p_c(t) = \frac{d}{dt} [w_c(t)]$
 \uparrow
 energy in the capacitor

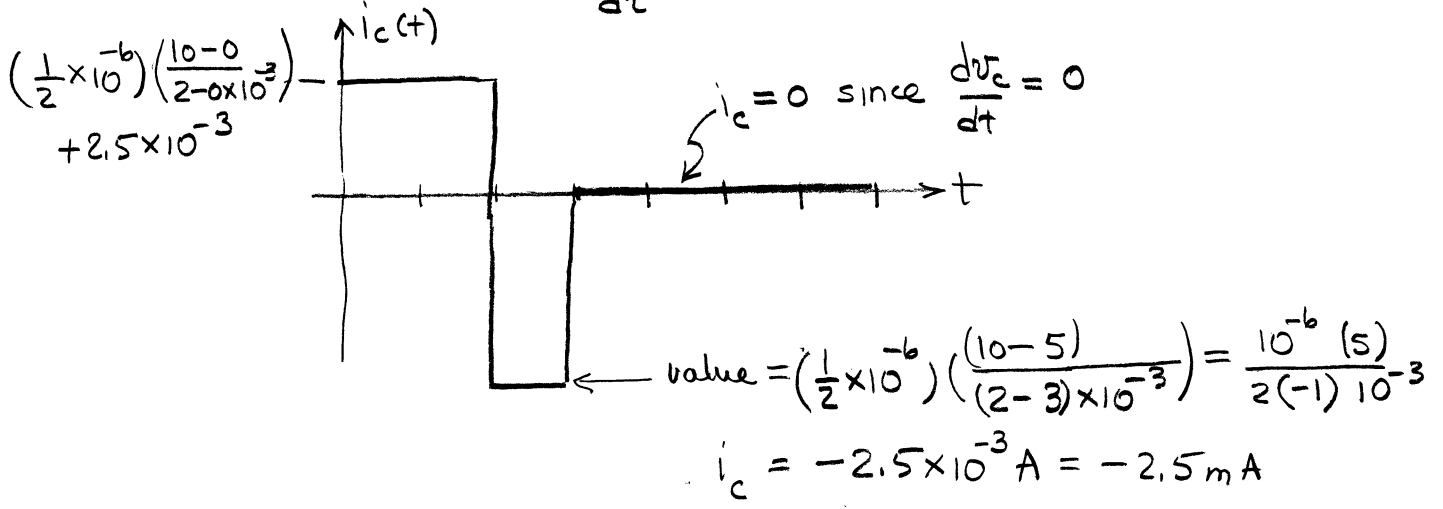
\therefore energy in capacitor $w_c(t) = \frac{1}{2} C v_c^2(t) + \text{constant}$
 since $\xi = 0$ when $v_c = 0$

Example 6-1

If the capacitor voltage is as shown across a $\frac{1}{2} \mu\text{F}$ capacitor. What is the corresponding current?



Solution: since $i_c = C \frac{dv_c}{dt}$



Example 6-2

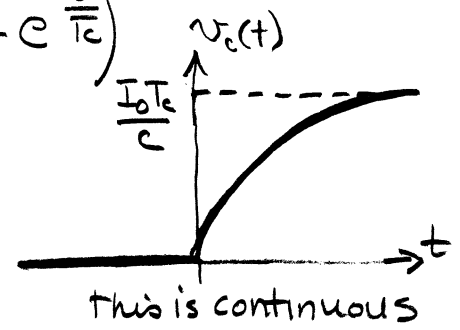
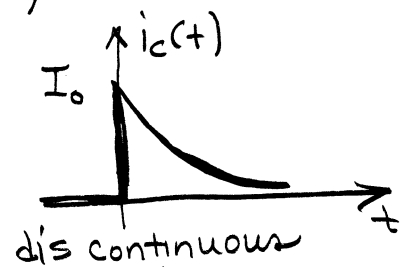
If $i_c(t) = I_0 e^{-\frac{t}{T_c}} u(t)$ find $v_c(t)$ if $v_c(t=0) = 0$.

$$i_c = C \frac{dv_c}{dt}$$

$$\frac{1}{C} i_c dt = dv_c$$

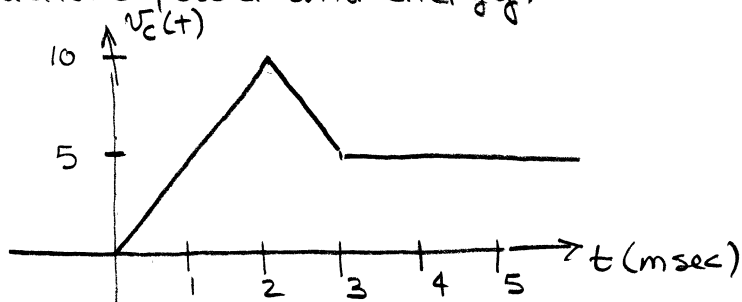
$$v_c(t) = v_c(t=0) + \frac{1}{C} \int_0^t I_0 e^{-\frac{x}{T_c}} dx$$

$$v_c(t) = 0 + \frac{I_0}{C} \left. \frac{e^{-\frac{x}{T_c}}}{-\frac{1}{T_c}} \right|_0^t = \frac{I_0 T_c}{C} (1 - e^{-\frac{t}{T_c}})$$



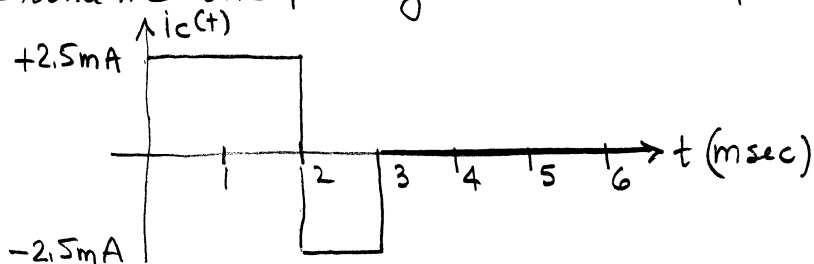
Example 6-3

The voltage across a capacitor is given below. Find the capacitor's power and energy.

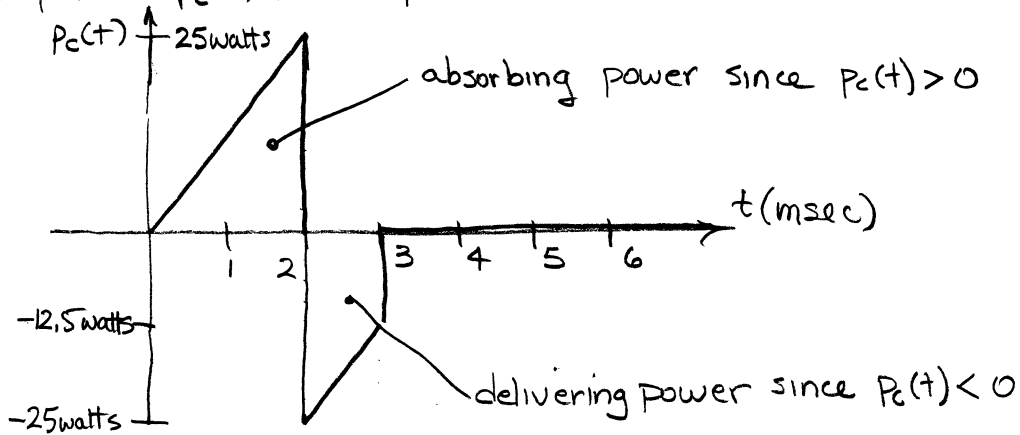


Solution:

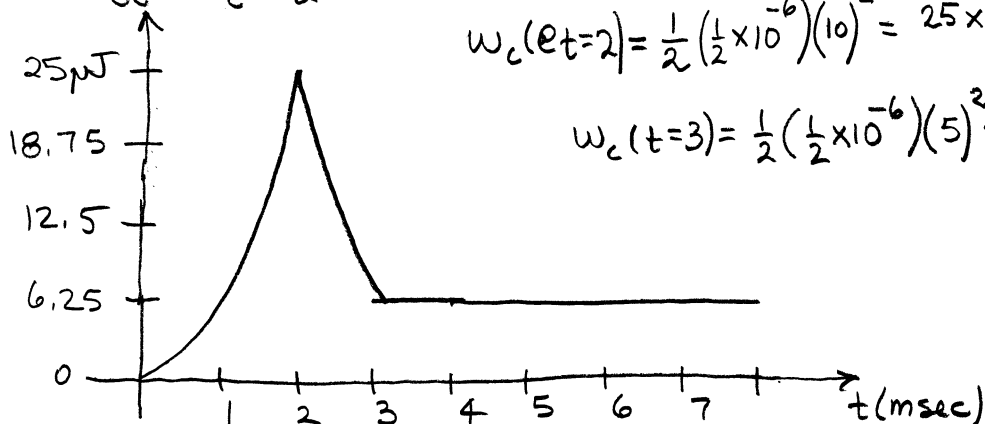
We found the corresponding current in Example 6-1.



The power $p_c(t)$ is the product of these waveforms.



The energy $w_c = \frac{1}{2} C V^2$



$$w_c(t=2) = \frac{1}{2} \left(\frac{1}{2} \times 10^{-6} \right) (10)^2 = 25 \times 10^{-6} \text{ J}$$

$$w_c(t=3) = \frac{1}{2} \left(\frac{1}{2} \times 10^{-6} \right) (5)^2 = 6.25 \times 10^{-6} \text{ J}$$

Example 6-4

The current through a capacitor is given by

$$i_c(t) = I_0 e^{-\frac{t}{\tau_c}} u(t)$$

Find the capacitor's energy and power.

The voltage was found in Example 6-2 to be

$$v_c(t) = \frac{I_0 \tau_c}{c} \left[1 - e^{-\frac{t}{\tau_c}} \right]$$

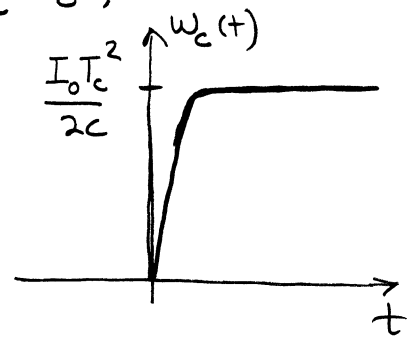
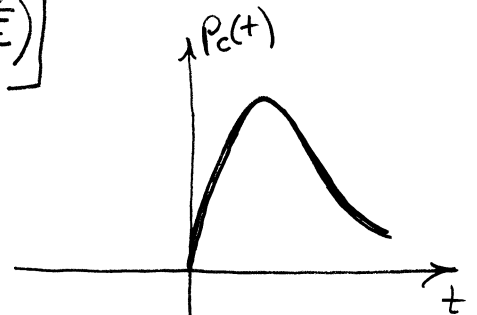
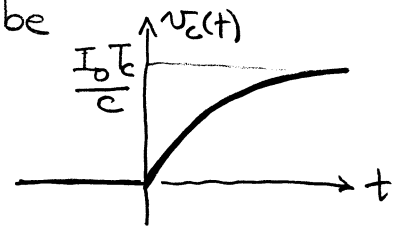
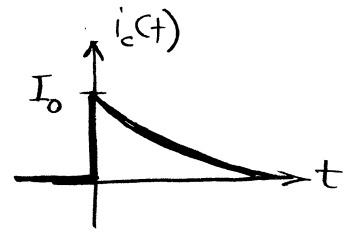
The power is given by $P_c(t) = i_c(t) v_c(t)$

$$P_c(t) = \left[I_0 e^{-\frac{t}{\tau_c}} \right] \left[\frac{I_0 \tau_c}{c} \left(1 - e^{-\frac{t}{\tau_c}} \right) \right]$$

$$P_c(t) = \frac{I_0^2 \tau_c}{c} \left[e^{-\frac{t}{\tau_c}} - e^{-\frac{2t}{\tau_c}} \right]$$

The energy is most easily calculated as $w_c(t) = \frac{1}{2} C v_c^2(t)$

$$w_c(t) = \frac{1}{2} C \left[\frac{I_0 \tau_c}{c} \left(1 - e^{-\frac{t}{\tau_c}} \right) \right]^2$$

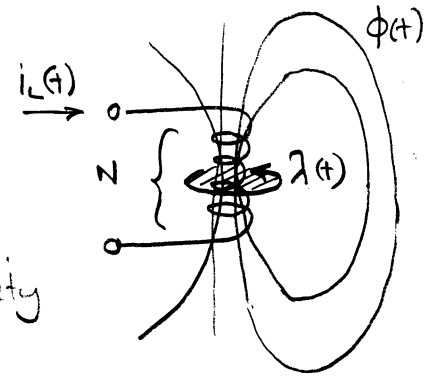


6-2 The Inductor

From magneto statics

$$\phi(t) = k_1 N i_L(t)$$

$\phi(t)$ magnetic flux (webers)
 k_1 constant of proportionality
 N # of turns
 $i_L(t)$ current



$$\lambda(t) = N \phi(t)$$

$\lambda(t)$ flux linkage, flux/unit area.
 N # of turns
 $\phi(t)$ magnetic flux

Substituting

$$\lambda(t) = N \phi(t) = N k_1 N i_L(t) = [N^2 k_1] i_L(t)$$

↑
this is the inductance L

$$\lambda(t) = L i_L(t)$$

Differentiate this to get the i-v relationship

$$\frac{d\lambda(t)}{dt} = L \frac{di_L(t)}{dt}$$

↑ by Faraday's Law $v_L(t) = \frac{d\lambda(t)}{dt}$

$$\therefore v_L(t) = L \frac{di_L(t)}{dt}$$

Observations

- ① If $i_L(t)$ is constant, $v_L = 0$ and the inductor looks like a short
- ② A discontinuity $w(t)$ or $\delta(t)$ in i_L would create an infinite voltage so $i_L(t)$ must be continuous.

We often want i_L in terms of voltage

$$v_L(t) = L \frac{di_L(t)}{dt}$$

$$di_L(t) = \frac{1}{L} v_L(t) dt$$

Assuming $i_L(t_0)$ is known we can integrate this to get

$$\int_{i_L(t_0)}^{i_L(t)} di_L = \frac{1}{L} \int_{t_0}^t v_L(x) dx$$

$$i_L(t) - i_L(t_0) = \frac{1}{L} \int_{t_0}^t v_L(x) dx$$

$$i_L(t) = i_L(t_0) + \frac{1}{L} \int_{t_0}^t v_L(x) dx$$

Power & energy

$$P_L(t) = i_L(t) v_L(t)$$

$$P_L(t) = i_L(t) L \frac{di_L(t)}{dt}$$

$$P_L(t) = \frac{d}{dt} \left[\frac{1}{2} L i_L^2(t) \right]$$

recognize this as the energy w_L
stored in the inductor

$$w_L(t) = \frac{1}{2} L i_L^2(t) + \text{constant}^0$$

since it is the energy
stored when $i_L = 0$

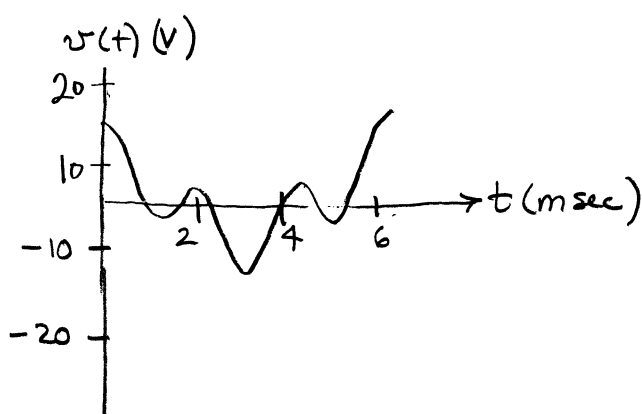
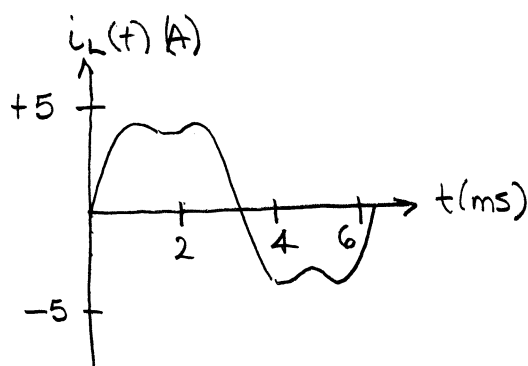
Example 6-6

The current through a 2-mH inductor is $i_L(t) = 4 \sin 1000t + \sin 3000t$. Find the corresponding $v_L(t)$.

$$v_L(t) = L \frac{di_L}{dt} = L \frac{d}{dt} (4 \sin 1000t + \sin 3000t)$$

$$= (.002) [4(1000) \cos 1000t + 3000 \cos 3000t]$$

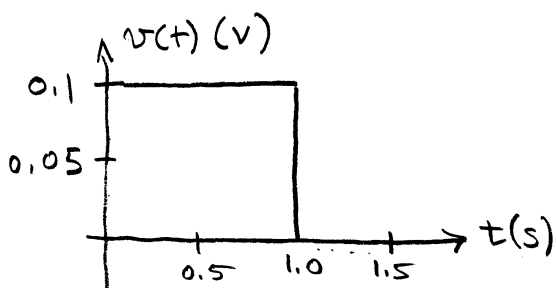
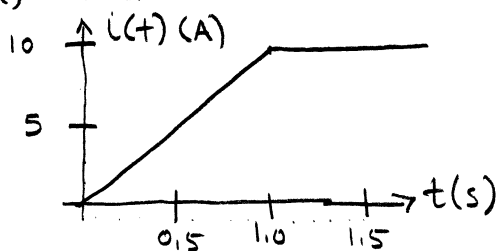
$$v_L(t) = 8 \cos 1000t + 6 \cos 3000t$$



Example 6-7

The following figures show the current and voltage across an unknown energy storage element.

(a) What is the element and its numerical value?



This is an inductor since $v=0$ when i is constant.

Since $v = L \frac{di}{dt}$ for $0 < t < 1$

$$0.1 = L \left(\frac{10-0}{1-0} \right) = 10L$$

$$L = .01 \text{ H} = 10 \text{ mH}$$

(b) What is the energy stored at $t=1 \text{ sec}$?

$$w_L = \frac{1}{2} L i^2 = \frac{1}{2} (.01) (10)^2 = \frac{1}{2} \text{ J}$$