

## Maxwell's Equations (general differential)

$$\nabla \times \mathbf{E} = -\frac{\partial \mathbf{B}}{\partial t} \quad [1.11a]$$

$$\nabla \cdot \mathbf{D} = \tilde{\rho} \quad [1.11b]$$

$$\nabla \times \mathbf{H} = \mathbf{J} + \frac{\partial \mathbf{D}}{\partial t} \quad [1.11c]$$

$$\nabla \cdot \mathbf{B} = 0 \quad [1.11d]$$

## Maxwell's Equations (time harmonic)

$$\nabla \times \mathbf{E} = -j\omega \mathbf{B} \quad [1.11a]$$

$$\nabla \cdot \mathbf{D} = \rho \quad [1.11b]$$

$$\nabla \times \mathbf{H} = \mathbf{J} + \frac{\partial \mathbf{D}}{\partial t} \quad [1.11c]$$

$$\nabla \cdot \mathbf{B} = 0 \quad [1.11d]$$

## Maxwell's Equations (integral)

$$\oint_C \mathbf{E} \cdot d\mathbf{l} = -\int_S \frac{\partial \mathbf{B}}{\partial t} \cdot d\mathbf{s} \quad [1.1]$$

$$\oint_S \mathbf{D} \cdot d\mathbf{s} = \int_V \rho dv \quad [1.2]$$

$$\oint_C \mathbf{H} \cdot d\mathbf{l} = \int_S \mathbf{J} \cdot d\mathbf{s} + \int_S \frac{\partial \mathbf{D}}{\partial t} \cdot d\mathbf{s} \quad [1.3]$$

$$\oint_S \mathbf{B} \cdot d\mathbf{s} = 0 \quad [1.4]$$

## Electromagnetic Boundary Conditions

$$\hat{\mathbf{n}} \times [\mathbf{E}_1 - \mathbf{E}_2] = 0 \quad [1.12]$$

$$\hat{\mathbf{n}} \times [\mathbf{H}_1 - \mathbf{H}_2] = 0 \quad [1.13]$$

$$\hat{\mathbf{n}} \times \mathbf{H}_1 = \mathbf{J} \quad [1.14]$$

$$\hat{\mathbf{n}} \cdot [\mathbf{D}_1 - \mathbf{D}_2] = \tilde{\rho}_s \quad [1.15]$$

$$\hat{\mathbf{n}} \cdot [\mathbf{B}_1 - \mathbf{B}_2] = 0 \quad [1.16]$$

$$\hat{\mathbf{n}} \cdot [\mathbf{J}_1 - \mathbf{J}_2] = -\frac{\partial \tilde{\rho}_s}{\partial t} \quad [1.17]$$

## Reflection &amp; transmission (simple dielectric)

$$\Gamma = \frac{E_{r0}}{E_{i0}} = \frac{\eta_2 - \eta_1}{\eta_2 + \eta_1} = \frac{n_1 - n_2}{n_1 + n_2} \quad [3.2]$$

$$T = \frac{E_{t0}}{E_{i0}} = \frac{2\eta_2}{\eta_2 + \eta_1} = \frac{2n_1}{n_1 + n_2} \quad [3.3]$$

## Basic waves

$$n \equiv \frac{c}{v_p} = \beta \frac{c}{\omega} = \sqrt{\mu_r \epsilon_r} = \sqrt{\frac{\epsilon}{\epsilon_0}} \quad \text{p.135}$$

$$\beta = \frac{2\pi}{\lambda} = \frac{2\pi f}{v_p} = \omega \sqrt{\mu \epsilon}$$

$$f\lambda = v_p = \frac{1}{\sqrt{\mu \epsilon}}; \quad \eta = \sqrt{\frac{\mu}{\epsilon}}$$

## Uniform plane waves in arbitrary direction

$$\underline{H}(\mathbf{r}) = \frac{1}{\eta} \hat{\mathbf{k}} \times \underline{E}(\mathbf{r}) \quad [2.61]$$

$$\underline{E}(\mathbf{r}) = -\eta \hat{\mathbf{k}} \times \underline{H}(\mathbf{r})$$

## Reflection &amp; transmission (multiple dielectrics)

$$\Gamma_{eff} = \frac{(\eta_2 - \eta_1)(\eta_3 + \eta_2) + (\eta_2 + \eta_1)(\eta_3 - \eta_2)e^{-2j\beta_2 d}}{(\eta_2 + \eta_1)(\eta_3 + \eta_2) + (\eta_2 - \eta_1)(\eta_3 - \eta_2)e^{-2j\beta_2 d}} \quad [3.7]$$

$$T_{eff} = \frac{4\eta_2 \eta_3 e^{-j\beta_2 d}}{(\eta_2 + \eta_1)(\eta_3 + \eta_2) + (\eta_2 - \eta_1)(\eta_3 - \eta_2)e^{-2j\beta_2 d}} \quad [3.8]$$

$$Z_2(-d) = \eta_2 \frac{\eta_3 + j\eta_2 \tan(\beta_2 d)}{\eta_2 + j\eta_3 \tan(\beta_2 d)} = \eta_{23} \quad [3.9]$$

$$\Gamma_{eff} = \rho_{eff} e^{j\phi_r} = \frac{Z_2(-d) - \eta_1}{Z_2(-d) + \eta_1} = \frac{\eta_{23} - \eta_1}{\eta_{23} + \eta_1} \quad [3.10]$$

## Plane waves in lossy materials

$$\gamma = \alpha + j\beta = \sqrt{j\omega\mu(\sigma + j\omega\epsilon)} \quad [2.19]$$

$$\alpha = \omega\sqrt{\frac{\mu\epsilon}{2} \left[ \sqrt{1 + \left(\frac{\sigma}{\omega\epsilon}\right)^2} - 1 \right]} \quad [2.20]$$

$$\beta = \omega\sqrt{\frac{\mu\epsilon}{2} \left[ \sqrt{1 + \left(\frac{\sigma}{\omega\epsilon}\right)^2} + 1 \right]} \quad [2.21]$$

## Lossy materials

Polarization currents:

$$\epsilon_c = \epsilon' - j\epsilon'' \quad [p.46]$$

$$\tan\delta_c = \frac{\sigma_{eff}}{\omega\epsilon'} = \frac{\epsilon''}{\epsilon'} \quad [p.48]$$

Conduction Currents:

$$\sigma_{eff} = \sigma + \omega\epsilon'' \quad [p.48]$$

$$\tan\delta_c = \frac{\sigma + \omega\epsilon''}{\omega\epsilon'} \quad [p.48]$$

Use  $\tan\delta_c$  instead of  $\frac{\sigma}{\omega\epsilon}$  in expression for complex impedance

$$\eta_c = \sqrt{\frac{\mu}{\epsilon_{eff}}} = \sqrt{\frac{\mu}{\epsilon - j\frac{\sigma}{\omega}}} = \frac{\sqrt{\frac{\mu}{\epsilon}}}{\left(1 + \left(\frac{\sigma}{\omega\epsilon}\right)^2\right)^{1/4}} e^{j(\frac{1}{2})\tan^{-1}(\frac{\sigma}{\omega\epsilon})} \quad [2.22]$$

## Good conductor approximations

$$\alpha = \beta \approx \sqrt{\frac{\omega\mu\sigma}{2}} \quad [p.54]$$

$$\delta \approx \frac{1}{\alpha} \approx \sqrt{\frac{2}{\omega\mu\sigma}} = \frac{1}{\sqrt{j\omega\mu\sigma}} \quad [2.26]$$

$$\eta_c \approx \sqrt{\frac{\mu}{\epsilon} \frac{1}{\left(1 - j\frac{\sigma}{\omega\epsilon}\right)^{1/2}}} = \sqrt{\frac{j\omega\mu}{\sigma}} = \sqrt{\frac{\omega\mu}{\sigma}} e^{j45^\circ} \quad [p.54]$$

## Poynting Theorem

$$\int_V \mathbf{E} \cdot \mathbf{J} dv = -\frac{\partial}{\partial t} \int_V \left( \frac{1}{2} \mu |\mathbf{H}|^2 + \frac{1}{2} \epsilon |\mathbf{E}|^2 \right) dv - \oint_S (\mathbf{E} \times \mathbf{H}) \cdot d\mathbf{s} \quad [2.31]$$

$$\mathbf{S}_{av} = \hat{z} \frac{E_0^2}{2\eta} \quad [2.32]$$

$$\mathbf{S}_{av} = \frac{1}{2} \mathbf{R} \{ \mathbf{E} \times \mathbf{H}^* \} \quad [2.43]$$

## Arbitrarily directed uniform plane waves

$$\mathbf{E}(\mathbf{r}) = \mathbf{E}_0 e^{-j\mathbf{k}\cdot\mathbf{r}} \quad [2.58]$$

$$\hat{\mathbf{k}} = \frac{\hat{\mathbf{x}}\beta_x + \hat{\mathbf{y}}\beta_y + \hat{\mathbf{z}}\beta_z}{\omega\sqrt{\mu\epsilon}} \quad [p.97]$$

$$\mathbf{H}(\mathbf{r}) = \frac{1}{\eta} \hat{\mathbf{k}} \times \mathbf{E}(\mathbf{r}) \quad [2.61]$$

## Reflection and refraction of oblique waves at planar dielectric interfaces

$$\frac{\sin\theta_i}{\sin\theta_t} = \frac{v_{p1}}{v_{p2}} = \frac{\sqrt{\epsilon_2\mu_2}}{\sqrt{\epsilon_1\mu_1}} = \frac{n_2}{n_1} \quad (\text{known as Snell's Law}) \quad [3.19]$$

$$\Gamma_{\perp} = \frac{E_{r0}}{E_{i0}} = \frac{\eta_2 \cos\theta_i - \eta_1 \cos\theta_t}{\eta_2 \cos\theta_i + \eta_1 \cos\theta_t} = \frac{\cos\theta_i - \sqrt{\frac{\epsilon_2}{\epsilon_1}} \sqrt{1 - \sin^2\theta_i}}{\cos\theta_i + \sqrt{\frac{\epsilon_2}{\epsilon_1}} \sqrt{1 - \sin^2\theta_i}} \quad [3.24]$$

$$T_{\perp} = \frac{E_{t0}}{E_{i0}} = \frac{2\eta_2 \cos\theta_i}{\eta_2 \cos\theta_i + \eta_1 \cos\theta_t} = \frac{2\cos\theta_i}{\cos\theta_i + \sqrt{\frac{\epsilon_2}{\epsilon_1}} \sqrt{1 - \sin^2\theta_i}} \quad [3.25]$$

$$\Gamma_{\parallel} = \frac{E_{r0}}{E_{i0}} = \frac{-\eta_1 \cos\theta_i + \eta_2 \cos\theta_t}{\eta_1 \cos\theta_i + \eta_2 \cos\theta_t} = \frac{-\cos\theta_i + \frac{\epsilon_1}{\epsilon_2} \sqrt{\frac{\epsilon_2}{\epsilon_1}} \sqrt{1 - \sin^2\theta_i}}{\cos\theta_i + \frac{\epsilon_1}{\epsilon_2} \sqrt{\frac{\epsilon_2}{\epsilon_1}} \sqrt{1 - \sin^2\theta_i}} \quad [3.26]$$

$$T_{\parallel} = \frac{E_{t0}}{E_{i0}} = \frac{2\eta_2 \cos\theta_i}{\eta_1 \cos\theta_i + \eta_2 \cos\theta_t} = \frac{2\sqrt{\frac{\epsilon_1}{\epsilon_2}} \cos\theta_i}{\cos\theta_i + \sqrt{\frac{\epsilon_1}{\epsilon_2}} \sqrt{1 - \frac{\epsilon_1}{\epsilon_2} \sin^2\theta_i}} \quad [3.27]$$

## Total internal reflection

$$\sin \theta_{ic} = \sqrt{\frac{\epsilon_2}{\epsilon_1}} = \frac{n_2}{n_1} \quad [3.29]$$

$$\text{for } \theta_i > \theta_{ic} \quad \cos \theta_t = \pm j \sqrt{\frac{\epsilon_1}{\epsilon_2} \sin^2 \theta_i - 1} \quad [3.31]$$

$$\Gamma_{\perp} = \frac{\cos \theta_i + j \sqrt{\sin^2 \theta_i - \frac{\epsilon_2}{\epsilon_1}}}{\cos \theta_i - j \sqrt{\sin^2 \theta_i - \frac{\epsilon_2}{\epsilon_1}}} = 1e^{j\phi_{\perp}} \quad [3.32]$$

$$\tan \frac{\phi_{\perp}}{2} = \frac{\sqrt{\sin^2 \theta_i - \frac{\epsilon_2}{\epsilon_1}}}{\cos \theta_i} \quad [\text{p.192}]$$

$$\Gamma_{\parallel} = -\frac{\frac{\epsilon_2}{\epsilon_1} \cos \theta_i + j \sqrt{\sin^2 \theta_i - \frac{\epsilon_2}{\epsilon_1}}}{\frac{\epsilon_2}{\epsilon_1} \cos \theta_i - j \sqrt{\sin^2 \theta_i - \frac{\epsilon_2}{\epsilon_1}}} = 1e^{j\phi_{\parallel}} \quad [3.33]$$

$$\tan \frac{\phi_{\parallel}}{2} = \frac{\sqrt{\sin^2 \theta_i - \frac{\epsilon_2}{\epsilon_1}}}{\frac{\epsilon_2}{\epsilon_1} \cos \theta_i} \quad [\text{p.192}]$$

## Normal incidence on a lossy medium

$$\gamma_2 = j\omega \sqrt{\mu \epsilon_{\text{eff}}} \approx j\omega \left( \frac{\mu_2 \sigma_2}{j\omega} \right)^{1/2} = (j\omega \mu_2 \sigma_2)^{1/2} = \frac{1+j}{\delta} \quad [3.39]$$

$$\eta_c = Z_s = R_s + jX_s = \sqrt{\frac{\mu_2}{\epsilon_{\text{eff}}}} \approx \left( \frac{j\omega \mu_2}{\sigma_2} \right)^{1/2} = \frac{\gamma_2}{\sigma_2} = \frac{1+j}{\sigma_2 \delta} \quad [3.40]$$

$$\Gamma_{\parallel} = -\frac{\frac{\epsilon_2}{\epsilon_1} \cos \theta_i + j \sqrt{\sin^2 \theta_i - \frac{\epsilon_2}{\epsilon_1}}}{\frac{\epsilon_2}{\epsilon_1} \cos \theta_i - j \sqrt{\sin^2 \theta_i - \frac{\epsilon_2}{\epsilon_1}}} = 1e^{j\phi_{\parallel}} \quad [3.33]$$

$$\tan \frac{\phi_{\parallel}}{2} = \frac{\sqrt{\sin^2 \theta_i - \frac{\epsilon_2}{\epsilon_1}}}{\frac{\epsilon_2}{\epsilon_1} \cos \theta_i} \quad [\text{p.192}]$$

## Parallel plate waveguide

Parallel-plate TE<sub>m</sub> modes: m=0,±1,±2,...

$$E_y = C_1 \sin\left(\frac{m\pi}{a} x\right) e^{-\tilde{\gamma} z} \quad [4.12a]$$

$$H_z = -\frac{1}{j\omega\mu} \frac{\partial E_y}{\partial x} = -\frac{m\pi}{j\omega\mu a} C_1 \cos\left(\frac{m\pi}{a} x\right) e^{-\tilde{\gamma} z} \quad [4.12b]$$

$$H_x = -\frac{\tilde{\gamma}}{j\omega\mu} E_y = -\frac{\tilde{\gamma}}{j\omega\mu} C_1 \sin\left(\frac{m\pi}{a} x\right) e^{-\tilde{\gamma} z} \quad [4.12c]$$

Parallel-plate TM<sub>m</sub> modes: m=0,±1,±2,...

$$H_y = C_4 \cos\left(\frac{m\pi}{a} x\right) e^{-\tilde{\gamma} z} \quad [4.13a]$$

$$E_x = \frac{\tilde{\gamma}}{j\omega\epsilon} H_y = \frac{\tilde{\gamma}}{j\omega\epsilon} C_4 \cos\left(\frac{m\pi}{a} x\right) e^{-\tilde{\gamma} z} \quad [4.13b]$$

$$E_z = \frac{j m \pi}{\omega \epsilon a} C_4 \sin\left(\frac{m\pi}{a} x\right) e^{-\tilde{\gamma} z} \quad [4.13c]$$

Parallel-plate TEM mode

$$H_y = C_4 e^{-\tilde{\gamma} z} \quad [4.14a]$$

$$E_x = \frac{\tilde{\gamma}}{j\omega\epsilon} C_4 e^{-\tilde{\gamma} z} \quad [4.14b]$$

$$E_z = 0 \quad [4.14c]$$

## Propagation constants

$$f_{cm} = \frac{mv_p}{2a} = \frac{m}{2a\sqrt{\mu\epsilon}} \quad [4.15]$$

$$\tilde{\gamma} = j\tilde{\beta}_m = j\sqrt{\omega^2 \mu \epsilon - \left(\frac{m\pi}{a}\right)^2} = j\beta \sqrt{1 - \left(\frac{f_{cm}}{f}\right)^2}, f > f_{cm} \quad [4.16]$$

$$\tilde{\gamma} = \tilde{\alpha}_m = j\sqrt{\left(\frac{m\pi}{a}\right)^2 - \omega^2 \mu \epsilon} = \beta \sqrt{\left(\frac{f_{cm}}{f}\right)^2 - 1}, f < f_{cm} \quad [4.14c]$$

$$\bar{\lambda}_m = \frac{2\pi}{\tilde{\beta}_m} = \frac{\lambda}{\sqrt{1 - \left(\frac{f_{cm}}{f}\right)^2}} \quad [4.18]$$

$$\bar{v}_{pm} = \frac{\omega}{\tilde{\beta}_m} = \frac{\bar{v}_p}{\sqrt{1 - \left(\frac{f_{cm}}{f}\right)^2}} \quad [4.18]$$

Conduction losses

$$\alpha_{cTM} = \frac{R_s}{\eta a} = \frac{1}{\eta a} \sqrt{\frac{\omega \mu_0}{2\sigma}} \quad [4.22]$$

$$\alpha_{cTE_m} = \frac{2R_s \left(\frac{f_{cm}}{f}\right)^2}{\eta a \sqrt{1 - \left(\frac{f_{cm}}{f}\right)^2}} \quad [4.23]$$

$$\alpha_{cTM_m} = \frac{2R_s}{\eta a \sqrt{1 - \left(\frac{f_{cm}}{f}\right)^2}} \quad [4.24]$$

Dielectric losses

$$\alpha_d = \frac{\omega \sqrt{\mu \epsilon'} \epsilon'' / \epsilon'^2}{2 \sqrt{1 - \left(\frac{\omega_{cm}}{\omega}\right)^2}} \quad [4.27]$$

For parallel plate TE modes the total power through the guide is

$$P_{av} = \int_0^b \int_0^a \frac{\tilde{B} C_1^2}{2\omega \mu} \sin^2\left(\frac{m\pi}{a}x\right) dx dy = \frac{\tilde{B} C_1^2 ab}{4\omega \mu} \quad [p.277]$$

For the parallel plate TEM (TM<sub>0</sub>) mode the total power through the guide is

$$P_{av} = \frac{1}{2} \eta C_4^2 ba \quad [p.277]$$

Dielectric slab waveguide

**TM Modes** The non-zero field components are  $E_z$ ,  $E_x$ , and  $H_y$

For  $x \leq -d/2$

$$E_z^0(x) = C_o \sin(\beta_x x) + C_e \cos(\beta_x x) \quad [4.34]$$

where the transverse propagation constant is given by

$$\beta_x^2 = \omega^2 \mu_d \epsilon_d - \beta^2 = h_d^2 \quad [4.35]$$

For  $x \geq d/2$

$$E_z^0(x) = \begin{cases} C_d e^{-\alpha_x(x-d/2)}, & x \geq d/2 \\ C_a e^{\alpha_x(x+d/2)}, & x \leq -d/2 \end{cases} \quad [4.36]$$

where the transverse attenuation constant is given by

$$\alpha_x^2 = \beta^2 - \omega^2 \mu_0 \epsilon_0 = -h_0^2 \text{ or } \alpha_x = \sqrt{\beta^2 - \omega^2 \mu_0 \epsilon_0} \quad [4.37]$$

The cutoff frequencies are given by

$$f_{cTMm} = \frac{(m-1)}{2d \sqrt{\mu_d \epsilon_d - \mu_0 \epsilon_0}} \quad m = 1, 3, 5, \dots \text{odd} \quad [4.45]$$

$$m = 2, 4, 6, \dots \text{even}$$

**TE Modes** The non-zero field components are  $H_z$ ,  $H_x$ , and  $E_y$

For  $x \leq -d/2$

$$H_z^0(x) = C_o \sin(\beta_x x) + C_e \cos(\beta_x x) \quad [4.46]$$

where the transverse propagation constant

$$\beta_x^2 = \omega^2 \mu_d \epsilon_d - \beta^2 = h_d^2$$

For  $x \geq d/2$

$$H_z^0(x) = \begin{cases} C_d e^{-\alpha_x(x-d/2)}, & x \geq d/2 \\ C_a e^{\alpha_x(x+d/2)}, & x \leq -d/2 \end{cases}$$

where the transverse propagation constant

$$\alpha_x^2 = \beta^2 - \omega^2 \mu_0 \epsilon_0 = -h_0^2 \text{ or } \alpha_x = \sqrt{\beta^2 - \omega^2 \mu_0 \epsilon_0} \quad [4.37]$$

For Odd TM Modes:

$$\frac{\alpha_x}{\beta_x} = \frac{\epsilon_0}{\epsilon_d} \tan\left(\frac{\beta_x d}{2}\right) \quad [4.40]$$

For Even TM Modes:

$$\frac{\alpha_x}{\beta_x} = -\frac{\epsilon_0}{\epsilon_d} \cot\left(\frac{\beta_x d}{2}\right) \quad [4.44]$$

For ALL TM Modes:

$$\alpha_x = \sqrt{\omega^2 (\mu_d \epsilon_d - \mu_0 \epsilon_0) - \beta_x^2} \quad [4.42]$$

The cutoff frequencies are given by

$$f_{cTMm} = \frac{(m-1)}{2d \sqrt{\mu_d \epsilon_d - \mu_0 \epsilon_0}} \quad m = 1, 3, 5, \dots \text{odd} \quad [4.49]$$

$$m = 2, 4, 6, \dots \text{even}$$

Dielectric covered ground plane

$$f_{c_{TE_m}} = \frac{(m-1)}{2d\sqrt{\mu_d\epsilon_d - \mu_0\epsilon_0}} \quad m = 1,3,5,\dots\text{odd\_}TM_m \quad [4.50]$$

Dielectric slab waveguide ray theory

$$\tan\theta_i = \frac{\bar{\beta}}{\beta_x} \quad [p.306]$$

Detailed example of odd TM Modes for slab dielectric waveguide: free space above the guide

For  $x \geq d/2$ 

$$E_z^0(x) = C_0 \sin\left(\frac{\beta_x d}{2}\right) e^{-\alpha_x(x-d/2)} \quad [4.39a]$$

$$E_x^0(x) = -\frac{j\bar{\beta}}{\alpha_x} C_0 \sin\left(\frac{\beta_x d}{2}\right) e^{-\alpha_x(x-d/2)} \quad [4.39b]$$

$$H_y^0(x) = -\frac{j\omega\epsilon_0}{\alpha_x} C_0 \sin\left(\frac{\beta_x d}{2}\right) e^{-\alpha_x(x-d/2)} \quad [4.39c]$$

For  $|x| \leq d/2$ 

$$E_z^0(x) = C_0 \sin(\beta_x x) \quad [4.39d]$$

$$E_x^0(x) = -\frac{j\bar{\beta}}{\beta_x} C_0 \cos(\beta_x x) \quad [4.39e]$$

$$H_y^0(x) = \frac{j\omega\epsilon_d}{\beta_x} C_0 \cos(\beta_x x) \quad [4.39f]$$

For  $x \leq -d/2$ 

$$E_z^0(x) = C_0 \sin\left(\frac{\beta_x d}{2}\right) e^{-\alpha_x(x+d/2)} \quad [4.39g]$$

$$E_x^0(x) = -\frac{j\bar{\beta}}{\alpha_x} C_0 \sin\left(\frac{\beta_x d}{2}\right) e^{-\alpha_x(x+d/2)} \quad [4.39h]$$

$$H_y^0(x) = -\frac{j\omega\epsilon_0}{\alpha_x} C_0 \sin\left(\frac{\beta_x d}{2}\right) e^{-\alpha_x(x+d/2)} \quad [4.39i]$$

Rectangular waveguides: TM modes

$$E_z = C \sin\left(\frac{m\pi}{a} x\right) \sin\left(\frac{n\pi}{b} y\right) e^{-j\bar{\beta}_{mn} z} \quad [5.13a]$$

$$E_x = -\frac{j\bar{\beta}_{mn} C}{h^2} \frac{m\pi}{a} \cos\left(\frac{m\pi}{a} x\right) \sin\left(\frac{n\pi}{b} y\right) e^{-j\bar{\beta}_{mn} z} \quad [5.13b]$$

$$E_y = -\frac{j\bar{\beta}_{mn} C}{h^2} \frac{n\pi}{b} \sin\left(\frac{m\pi}{a} x\right) \cos\left(\frac{n\pi}{b} y\right) e^{-j\bar{\beta}_{mn} z} \quad [5.13c]$$

$$H_x = \frac{j\omega\epsilon C}{h^2} \frac{n\pi}{b} \sin\left(\frac{m\pi}{a} x\right) \cos\left(\frac{n\pi}{b} y\right) e^{-j\bar{\beta}_{mn} z} \quad [5.13d]$$

$$H_y = -\frac{j\omega\epsilon C}{h^2} \frac{m\pi}{a} \cos\left(\frac{m\pi}{a} x\right) \sin\left(\frac{n\pi}{b} y\right) e^{-j\bar{\beta}_{mn} z} \quad [5.13e]$$

$$Z_{TM_{mn}} = \frac{E_x^0}{H_y^0} = \eta \sqrt{1 - \left(\frac{f_{c_{mn}}}{f}\right)^2} \quad [5.16]$$

Rectangular waveguides: TE modes

$$H_z = C \cos\left(\frac{m\pi}{a} x\right) \cos\left(\frac{n\pi}{b} y\right) e^{-j\bar{\beta}_{mn} z} \quad [5.21a]$$

$$H_x = \frac{j\bar{\beta}_{mn} C}{h^2} \frac{m\pi}{a} \sin\left(\frac{m\pi}{a} x\right) \cos\left(\frac{n\pi}{b} y\right) e^{-j\bar{\beta}_{mn} z} \quad [5.21b]$$

$$H_y = \frac{j\bar{\beta}_{mn} C}{h^2} \frac{n\pi}{b} \cos\left(\frac{m\pi}{a} x\right) \sin\left(\frac{n\pi}{b} y\right) e^{-j\bar{\beta}_{mn} z} \quad [5.21c]$$

$$E_x = \frac{j\omega\mu}{h^2} C \frac{n\pi}{b} \cos\left(\frac{m\pi}{a} x\right) \sin\left(\frac{n\pi}{b} y\right) e^{-j\bar{\beta}_{mn} z} \quad [5.21d]$$

$$E_y = -\frac{j\omega\mu}{h^2} C \frac{m\pi}{a} \sin\left(\frac{m\pi}{a} x\right) \cos\left(\frac{n\pi}{b} y\right) e^{-j\bar{\beta}_{mn} z} \quad [5.21e]$$

$$Z_{TE_{mn}} = \frac{\eta}{\sqrt{1 - \left(\frac{f_{c_{mn}}}{f}\right)^2}} \quad [5.22]$$

For both TM and TE modes:

$$\bar{\gamma} = \sqrt{h^2 - \omega^2 \mu\epsilon} = \sqrt{\left(\frac{m\pi}{a}\right)^2 + \left(\frac{n\pi}{b}\right)^2 - \omega^2 \mu\epsilon} \quad [5.14]$$

$$\text{where } h_{mn}^2 = \left(\frac{m\pi}{a}\right)^2 + \left(\frac{n\pi}{b}\right)^2$$

$$\omega_{c_{mn}} = \frac{1}{\sqrt{\mu\epsilon}} \sqrt{\left(\frac{m\pi}{a}\right)^2 + \left(\frac{n\pi}{b}\right)^2} \quad [5.15]$$

$$\bar{\beta}_{mn} = \sqrt{\omega^2 \mu\epsilon - \left(\frac{m\pi}{a}\right)^2 - \left(\frac{n\pi}{b}\right)^2} = \beta \sqrt{1 - \left(\frac{f_{c_{mn}}}{f}\right)^2} \quad [5.16]$$

Dominant TE<sub>10</sub> mode:

$$H_z = C \cos\left(\frac{\pi x}{a}\right) e^{-j\bar{\beta}_{10}z} \quad [5.23a]$$

$$H_x = \frac{j\bar{\beta}_{10}aC}{\pi} \sin\left(\frac{\pi x}{a}\right) e^{-j\bar{\beta}_{10}z} \quad [5.23b]$$

$$E_y = -\frac{j\omega\mu aC}{\pi} \sin\left(\frac{\pi x}{a}\right) e^{-j\bar{\beta}_{10}z} \quad [5.23c]$$

$$E_x = H_y = 0 \quad [5.23d]$$

$$\bar{\beta}_{10} = \sqrt{\omega^2\mu\epsilon - \left(\frac{\pi}{a}\right)^2} = \sqrt{\left(\frac{2\pi}{\lambda}\right)^2 - \left(\frac{\pi}{a}\right)^2} \quad [5.23e]$$

$$\alpha_{cTE_{10}} = \frac{R_s}{\eta b} \left[ \frac{1 + \frac{2b}{a} \left(\frac{f_{c10}}{f}\right)^2}{\sqrt{1 - \left(\frac{f_{c10}}{f}\right)^2}} \right] \quad [5.29]$$

$$\bar{\lambda}_{10} = \frac{2\pi}{\bar{\beta}_{10}} = \frac{\lambda}{\sqrt{1 - \left(\frac{\lambda}{2a}\right)^2}} \quad [5.23f]$$

$$f_{c10} = \frac{1}{2a\sqrt{\mu\epsilon}} \quad [5.24]$$

$$Z_{TE_{10}} = \frac{\eta}{\sqrt{1 - \left(\frac{f_{c10}}{f}\right)^2}} = \frac{2a\eta f}{\sqrt{4a^2\mu\epsilon f^2 - 1}} \quad [5.24]$$

Circular waveguides: TM modes where  $H_z = 0$ 

$$E_z = C_n J_n\left(\frac{t_{nl}}{a} r\right) \cos(n\phi) e^{-j\bar{\beta}_{nl}z} \quad [5.46a]$$

$$E_r = -\frac{j a \bar{\beta}_{TM_{nl}}}{t_{nl}} C_n J_n\left(\frac{t_{nl}}{a} r\right) \cos(n\phi) e^{-j\bar{\beta}_{nl}z} \quad [5.46b]$$

$$E_\phi = -\frac{j a^2 n \bar{\beta}_{TM_{nl}}}{t_{nl}^2} C_n J_n\left(\frac{t_{nl}}{a} r\right) \sin(n\phi) e^{-j\bar{\beta}_{nl}z} \quad [5.46c]$$

$$H_r = -\frac{j\omega\epsilon n a^2}{t_{nl}^2} C_n J_n\left(\frac{t_{nl}}{a} r\right) \sin(n\phi) e^{-j\bar{\beta}_{nl}z} \quad [5.46d]$$

$$H_\phi = -\frac{j\omega\epsilon a}{t_{nl}} C_n J_n\left(\frac{t_{nl}}{a} r\right) \cos(n\phi) e^{-j\bar{\beta}_{nl}z} \quad [5.46e]$$

where

$$\bar{\beta}_{TM_{nl}} = \sqrt{\omega^2\mu\epsilon - \left(\frac{t_{nl}}{a}\right)^2} \quad [5.44]$$

$$f_{cTM_{nl}} = \frac{t_{nl}}{2\pi a \sqrt{\mu\epsilon}} \quad [5.45]$$

Circular waveguides: TE modes where  $E_z = 0$ 

$$H_z = C_n J_n\left(\frac{s_{nl}}{a} r\right) \cos(n\phi) e^{-j\bar{\beta}_{nl}z} \quad [5.48a]$$

$$H_r = -\frac{j a \bar{\beta}_{TE_{nl}}}{s_{nl}} C_n J_n\left(\frac{s_{nl}}{a} r\right) \cos(n\phi) e^{-j\bar{\beta}_{nl}z} \quad [5.48b]$$

$$H_\phi = \frac{j n a^2 \bar{\beta}_{TE_{nl}}}{s_{nl}^2} C_n J_n\left(\frac{s_{nl}}{a} r\right) \sin(n\phi) e^{-j\bar{\beta}_{nl}z} \quad [5.48c]$$

$$E_r = -\frac{j a^2 \omega \mu n}{s_{nl}^2} C_n J_n\left(\frac{s_{nl}}{a} r\right) \sin(n\phi) e^{-j\bar{\beta}_{nl}z} \quad [5.48d]$$

$$E_\phi = -\frac{j a \omega \mu}{s_{nl}} C_n J_n\left(\frac{s_{nl}}{a} r\right) \cos(n\phi) e^{-j\bar{\beta}_{nl}z} \quad [5.48e]$$

where

$$\bar{\beta}_{TE_{nl}} = \sqrt{\omega^2\mu\epsilon - \left(\frac{s_{nl}}{a}\right)^2} \quad [p.359]$$

$$f_{cTE_{nl}} = \frac{s_{nl}}{2\pi a \sqrt{\mu\epsilon}} \quad [p.359]$$

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Table 5.2  $l$ th roots ( $t_{nl}$ ) of  $J_n(\cdot)=0$

$l=$	$n=$							
	0	1	2	3	4	5	6	7
1	2.405	3.832	5.136	6.380	7.588	8.771	9.936	11.086
2	5.520	7.016	8.417	9.761	11.065	12.339	13.589	14.821
3	8.654	10.173	11.620	13.015	14.372	15.700	17.004	18.288
4	11.792	13.323	14.796	16.223	17.616	18.980	20.321	21.642

Table 5.3  $l$ th roots ( $s_{nl}$ ) of  $J_n'(\cdot)=0$

$l=$	$n=$							
	0	1	2	3	4	5	6	7
1	3.832	1.841	3.054	4.201	5.317	6.416	7.501	8.578
2	7.016	5.331	6.706	8.015	9.282	10.520	11.735	12.932
3	10.173	8.536	9.969	11.346	12.682	13.987	15.268	16.529
4	13.324	11.706	13.170	14.586	15.964	17.313	18.637	19.942