

## Maxwell's Equations (general differential)

$$\nabla \times \vec{\mathcal{E}} = -\frac{\partial \vec{\mathcal{B}}}{\partial t} \quad [1.11a]$$

$$\nabla \cdot \vec{\mathcal{D}} = \tilde{\rho} \quad [1.11b]$$

$$\nabla \times \vec{\mathcal{H}} = \vec{\mathcal{J}} + \frac{\partial \vec{\mathcal{D}}}{\partial t} \quad [1.11c]$$

$$\nabla \cdot \vec{\mathcal{B}} = 0 \quad [1.11d]$$

## Maxwell's Equations (time harmonic)

$$\nabla \times \mathbf{E} = -j\omega \mathbf{B} \quad [1.11a]$$

$$\nabla \cdot \mathbf{D} = \rho \quad [1.11b]$$

$$\nabla \times \mathbf{H} = \mathbf{J} + \frac{\partial \mathbf{D}}{\partial t} \quad [1.11c]$$

$$\nabla \cdot \mathbf{B} = 0 \quad [1.11d]$$

## Maxwell's Equations (integral)

$$\oint_C \vec{\mathcal{E}} \cdot d\mathbf{l} = -\int_S \frac{\partial \vec{\mathcal{B}}}{\partial t} \cdot d\mathbf{s} \quad [1.1]$$

$$\oint_S \vec{\mathcal{D}} \cdot d\mathbf{s} = \int_V \rho dv \quad [1.2]$$

$$\oint_C \vec{\mathcal{H}} \cdot d\mathbf{l} = \int_S \vec{\mathcal{J}} \cdot d\mathbf{s} + \int_S \frac{\partial \vec{\mathcal{D}}}{\partial t} \cdot d\mathbf{s} \quad [1.3]$$

$$\oint_S \vec{\mathcal{B}} \cdot d\mathbf{s} = 0 \quad [1.4]$$

## Electromagnetic Boundary Conditions

$$\hat{\mathbf{n}} \times [\vec{\mathcal{E}}_1 - \vec{\mathcal{E}}_2] = 0 \quad [1.12]$$

$$\hat{\mathbf{n}} \times [\vec{\mathcal{H}}_1 - \vec{\mathcal{H}}_2] = 0 \quad [1.13]$$

$$\hat{\mathbf{n}} \times \vec{\mathcal{H}}_1 = \vec{\mathcal{J}} \quad [1.14]$$

$$\hat{\mathbf{n}} \cdot [\vec{\mathcal{D}}_1 - \vec{\mathcal{D}}_2] = \tilde{\rho}_s \quad [1.15]$$

$$\hat{\mathbf{n}} \cdot [\vec{\mathcal{B}}_1 - \vec{\mathcal{B}}_2] = 0 \quad [1.16]$$

$$\hat{\mathbf{n}} \cdot [\vec{\mathcal{J}}_1 - \vec{\mathcal{J}}_2] = -\frac{\partial \tilde{\rho}_s}{\partial t} \quad [1.17]$$

## Reflection &amp; transmission (simple dielectric)

$$\Gamma = \frac{E_{r0}}{E_{i0}} = \frac{\eta_2 - \eta_1}{\eta_2 + \eta_1} \quad [3.2]$$

$$T = \frac{E_{t0}}{E_{i0}} = \frac{2\eta_2}{\eta_2 + \eta_1} \quad [3.3]$$

## Basic waves

$$n \equiv \frac{c}{v_p} = \beta \frac{c}{\omega} = \sqrt{\mu_r \epsilon_r} \quad \text{p.135}$$

$$\beta = \frac{2\pi}{\lambda} = \frac{2\pi f}{v_p} = \omega \sqrt{\mu \epsilon}$$

$$f\lambda = v_p = \frac{1}{\sqrt{\mu \epsilon}}$$

## Reflection &amp; transmission (multiple dielectrics)

$$\Gamma_{eff} = \frac{(\eta_2 - \eta_1)(\eta_3 + \eta_2) + (\eta_2 + \eta_1)(\eta_3 - \eta_2)e^{-2j\beta_2 d}}{(\eta_2 + \eta_1)(\eta_3 + \eta_2) + (\eta_2 - \eta_1)(\eta_3 - \eta_2)e^{-2j\beta_2 d}} \quad [3.7]$$

$$T_{eff} = \frac{4\eta_2 \eta_3 e^{-j\beta_2 d}}{(\eta_2 + \eta_1)(\eta_3 + \eta_2) + (\eta_2 - \eta_1)(\eta_3 - \eta_2)e^{-2j\beta_2 d}} \quad [3.8]$$

$$Z_2(-d) = \eta_2 \frac{\eta_3 + j\eta_2 \tan(\beta_2 d)}{\eta_2 + j\eta_3 \tan(\beta_2 d)} = \eta_{23} \quad [3.9]$$

$$\Gamma_{eff} = \rho_{eff} e^{j\phi_r} = \frac{Z_2(-d) - \eta_1}{Z_2(-d) + \eta_1} = \frac{\eta_{23} - \eta_1}{\eta_{23} + \eta_1} \quad [3.10]$$

## Plane waves in lossy materials

$$\gamma = \alpha + j\beta = \sqrt{j\omega\mu(\sigma + j\omega\epsilon)} \quad [2.19]$$

$$\alpha = \omega \sqrt{\frac{\mu\epsilon}{2} \left[ \sqrt{1 + \left(\frac{\sigma}{\omega\epsilon}\right)^2} - 1 \right]} \quad [2.20]$$

$$\beta = \omega \sqrt{\frac{\mu\epsilon}{2} \left[ \sqrt{1 + \left(\frac{\sigma}{\omega\epsilon}\right)^2} + 1 \right]} \quad [2.21]$$

## Lossy materials

## Polarization currents:

$$\epsilon_c = \epsilon' - j\epsilon'' \quad [p.46]$$

$$\tan \delta_c = \frac{\sigma_{eff}}{\omega\epsilon'} = \frac{\epsilon''}{\epsilon'} \quad [p.48]$$

## Conduction currents:

$$\sigma_{eff} = \sigma + \omega\epsilon'' \quad [p.48]$$

$$\tan \delta_c = \frac{\sigma + \omega\epsilon''}{\omega\epsilon'} \quad [p.48]$$

Use  $\tan \delta_c$  instead of  $\frac{\sigma}{\omega\epsilon}$  in expression for complex impedance

$$\eta_c = \sqrt{\frac{\mu}{\epsilon_{eff}}} = \sqrt{\frac{\mu}{\epsilon - j\frac{\sigma}{\omega}}} = \frac{\sqrt{\frac{\mu}{\epsilon}}}{\left(1 + \left(\frac{\sigma}{\omega\epsilon}\right)^2\right)^{1/4}} e^{j(1/2)\tan^{-1}(\sigma/\omega\epsilon)} \quad [2.22]$$

## Good conductor approximations

$$\alpha = \beta \approx \sqrt{\frac{\omega\mu\sigma}{2}} \quad [p.54]$$

$$\delta \equiv \frac{1}{\alpha} \approx \sqrt{\frac{2}{\omega\mu\sigma}} = \frac{1}{\sqrt{\pi f \mu \sigma}} \quad [2.26]$$

$$\eta_c = \sqrt{\frac{\mu}{\epsilon} \frac{1}{\left(1 - j\frac{\sigma}{\omega\epsilon}\right)^{1/2}}} = \sqrt{\frac{j\omega\mu}{\sigma}} = \sqrt{\frac{\omega\mu}{\sigma}} e^{j45^\circ} \quad [p.54]$$

## Poynting Theorem

$$\int_V \bar{\mathcal{E}} \cdot \bar{\mathcal{J}} dv = -\frac{\partial}{\partial t} \int_V \left( \frac{1}{2} \mu |\bar{\mathcal{H}}|^2 + \frac{1}{2} \epsilon |\bar{\mathcal{E}}|^2 \right) dv - \oint_S (\bar{\mathcal{E}} \times \bar{\mathcal{H}}) \cdot d\mathbf{s} \quad [2.31]$$

$$\mathbf{S}_{av} = \hat{z} \frac{E_0^2}{2\eta} \quad [2.32]$$

$$\mathbf{S}_{av} = \frac{1}{2} \mathcal{R}_e \{ \mathbf{E} \times \mathbf{H}^* \} \quad [2.43]$$

## Arbitrarily directed uniform plane waves

$$\mathbf{E}(\mathbf{r}) = \mathbf{E}_0 e^{-j\mathbf{k} \cdot \mathbf{r}} \quad [2.58]$$

$$\hat{\mathbf{k}} = \frac{\hat{\mathbf{x}}\beta_x + \hat{\mathbf{y}}\beta_y + \hat{\mathbf{z}}\beta_z}{\omega\sqrt{\mu\epsilon}} \quad [p.97]$$

$$\mathbf{H}(\mathbf{r}) = \frac{1}{\eta} \hat{\mathbf{k}} \times \mathbf{E}(\mathbf{r}) \quad [2.61]$$

## Reflection and refraction of oblique waves at planar dielectric interfaces

$$\frac{\sin\theta_i}{\sin\theta_t} = \frac{v_{p1}}{v_{p2}} = \sqrt{\frac{\epsilon_2\mu_2}{\epsilon_1\mu_1}} \quad [3.19]$$

$$\Gamma_{\perp} = \frac{E_{r0}}{E_{i0}} = \frac{\eta_2 \cos\theta_i - \eta_1 \cos\theta_t}{\eta_2 \cos\theta_i + \eta_1 \cos\theta_t} = \frac{\cos\theta_i - \sqrt{\frac{\epsilon_2}{\epsilon_1} - \sin^2\theta_i}}{\cos\theta_i + \sqrt{\frac{\epsilon_2}{\epsilon_1} - \sin^2\theta_i}} \quad [3.24]$$

$$T_{\perp} = \frac{E_{t0}}{E_{i0}} = \frac{2\eta_2 \cos\theta_i}{\eta_2 \cos\theta_i + \eta_1 \cos\theta_t} = \frac{2\cos\theta_i}{\cos\theta_i + \sqrt{\frac{\epsilon_2}{\epsilon_1} - \sin^2\theta_i}} \quad [3.25]$$

$$\Gamma_{\parallel} = \frac{E_{r0}}{E_{i0}} = \frac{-\eta_1 \cos\theta_i + \eta_2 \cos\theta_t}{\eta_1 \cos\theta_i + \eta_2 \cos\theta_t} = \frac{-\cos\theta_i + \frac{\epsilon_1}{\epsilon_2} \sqrt{\frac{\epsilon_2}{\epsilon_1} - \sin^2\theta_i}}{\cos\theta_i + \frac{\epsilon_1}{\epsilon_2} \sqrt{\frac{\epsilon_2}{\epsilon_1} - \sin^2\theta_i}} \quad [3.26]$$

$$T_{\parallel} = \frac{E_{t0}}{E_{i0}} = \frac{2\eta_2 \cos\theta_i}{\eta_1 \cos\theta_i + \eta_2 \cos\theta_t} = \frac{2\sqrt{\frac{\epsilon_1}{\epsilon_2}} \cos\theta_i}{\cos\theta_i + \sqrt{\frac{\epsilon_1}{\epsilon_2} \left( 1 - \frac{\epsilon_1}{\epsilon_2} \sin^2\theta_i \right)}} \quad [3.27]$$

## Total internal reflection

$$\Gamma_{\perp} = \frac{\cos\theta_i + j\sqrt{\sin^2\theta_i - \frac{\epsilon_2}{\epsilon_1}}}{\cos\theta_i - j\sqrt{\sin^2\theta_i - \frac{\epsilon_2}{\epsilon_1}}} = 1e^{j\phi_r} \quad [3.32]$$

$$\tan \frac{\phi_r}{2} = \frac{\sqrt{\sin^2\theta_i - \frac{\epsilon_2}{\epsilon_1}}}{\cos\theta_i} \quad [p.192]$$

$$\Gamma_{\parallel} = -\frac{\frac{\epsilon_2}{\epsilon_1} \cos\theta_i + j\sqrt{\sin^2\theta_i - \frac{\epsilon_2}{\epsilon_1}}}{\frac{\epsilon_2}{\epsilon_1} \cos\theta_i - j\sqrt{\sin^2\theta_i - \frac{\epsilon_2}{\epsilon_1}}} = 1e^{j\phi_r} \quad [3.33]$$

$$\tan \frac{\phi_r}{2} = \frac{\sqrt{\sin^2\theta_i - \frac{\epsilon_2}{\epsilon_1}}}{\frac{\epsilon_2}{\epsilon_1} \cos\theta_i} \quad [p.192]$$

## Normal incidence on a lossy medium

$$\gamma_2 = j\omega\sqrt{\mu\epsilon_{eff}} \approx j\omega\left(\frac{\mu_2\sigma_2}{j\omega}\right)^{1/2} = (j\omega\mu_2\sigma_2)^{1/2} = \frac{1+j}{\delta} \quad [3.39]$$

$$\eta_c = Z_s = R_s + jX_s = \sqrt{\frac{\mu_2}{\epsilon_{eff}}} \approx \left(\frac{j\omega\mu_2}{\sigma_2}\right)^{1/2} = \frac{\gamma_2}{\sigma_2} = \frac{1+j}{\sigma_2\delta} \quad [3.40]$$

$$\Gamma_{\parallel} = -\frac{\frac{\epsilon_2}{\epsilon_1} \cos\theta_i + j\sqrt{\sin^2\theta_i - \frac{\epsilon_2}{\epsilon_1}}}{\frac{\epsilon_2}{\epsilon_1} \cos\theta_i - j\sqrt{\sin^2\theta_i - \frac{\epsilon_2}{\epsilon_1}}} = 1e^{j\phi_r} \quad [3.33]$$

$$\tan \frac{\phi_r}{2} = \frac{\sqrt{\sin^2\theta_i - \frac{\epsilon_2}{\epsilon_1}}}{\frac{\epsilon_2}{\epsilon_1} \cos\theta_i} \quad [p.192]$$

Parallel plate slab waveguide

 Parallel-plate TE<sub>m</sub> modes: m=0,±1,±2,...

$$E_y = C_1 \sin\left(\frac{m\pi}{a}x\right)e^{-\tilde{\gamma}z} \quad [4.12a]$$

$$H_z = -\frac{1}{j\omega\mu} \frac{\partial E_y}{\partial x} = -\frac{m\pi}{j\omega\mu a} C_1 \cos\left(\frac{m\pi}{a}x\right)e^{-\tilde{\gamma}z} \quad [4.12b]$$

$$H_x = -\frac{\tilde{\gamma}}{j\omega\mu} E_y = -\frac{\tilde{\gamma}}{j\omega\mu} C_1 \sin\left(\frac{m\pi}{a}x\right)e^{-\tilde{\gamma}z} \quad [4.12c]$$

 Parallel-plate TM<sub>m</sub> modes: m=0,±1,±2,...

$$H_y = C_4 \cos\left(\frac{m\pi}{a}x\right)e^{-\tilde{\gamma}z} \quad [4.13a]$$

$$E_x = \frac{\tilde{\gamma}}{j\omega\epsilon} H_y = \frac{\tilde{\gamma}}{j\omega\epsilon} C_4 \cos\left(\frac{m\pi}{a}x\right)e^{-\tilde{\gamma}z} \quad [4.13b]$$

$$E_z = \frac{j m \pi}{\omega \epsilon a} C_4 \sin\left(\frac{m\pi}{a}x\right)e^{-\tilde{\gamma}z} \quad [4.13c]$$

Parallel-plate TEM mode

$$H_y = C_4 e^{-\tilde{\gamma}z} \quad [4.14a]$$

$$E_x = \frac{\tilde{\gamma}}{j\omega\epsilon} C_4 e^{-\tilde{\gamma}z} \quad [4.14b]$$

$$E_z = 0 \quad [4.14c]$$

Propagation constants

$$f_{c_m} = \frac{m v_p}{2a} = \frac{m}{2a\sqrt{\mu\epsilon}} \quad [4.15]$$

$$\tilde{\gamma} = j\tilde{\beta}_m = j\sqrt{\omega^2\mu\epsilon - \left(\frac{m\pi}{a}\right)^2} = j\beta\sqrt{1 - \left(\frac{f_{c_m}}{f}\right)^2}, f > f_{c_m} \quad [4.16]$$

$$\tilde{\gamma} = \tilde{\alpha}_m = j\sqrt{\left(\frac{m\pi}{a}\right)^2 - \omega^2\mu\epsilon} = \beta\sqrt{\left(\frac{f_{c_m}}{f}\right)^2 - 1}, f < f_{c_m} \quad [4.14c]$$

$$\tilde{\lambda}_m = \frac{2\pi}{\tilde{\beta}_m} = \frac{\lambda}{\sqrt{1 - \left(\frac{f_{c_m}}{f}\right)^2}} \quad [4.18]$$

$$\bar{v}_{p_m} = \frac{\omega}{\tilde{\beta}_m} = \frac{\bar{v}_p}{\sqrt{1 - \left(\frac{f_{c_m}}{f}\right)^2}} \quad [4.18]$$

Conduction losses

$$\alpha_{c_{TEM}} = \frac{R_s}{\eta a} = \frac{1}{\eta a} \sqrt{\frac{\omega\mu_0}{2\sigma}} \quad [4.22]$$

$$\alpha_{c_{TE_m}} = \frac{2R_s \left(\frac{f_{c_m}}{f}\right)^2}{\eta a \sqrt{1 - \left(\frac{f_{c_m}}{f}\right)^2}} \quad [4.23]$$

$$\alpha_{c_{TM_m}} = \frac{2R_s}{\eta a \sqrt{1 - \left(\frac{f_{c_m}}{f}\right)^2}} \quad [4.24]$$

Dielectric losses

$$\alpha_d = \frac{\omega \sqrt{\mu\epsilon'} \epsilon'' / \epsilon'}{2 \sqrt{1 - \left(\frac{f_{c_m}}{f}\right)^2}} \quad [4.27]$$

For parallel plate TE modes the total power through the guide is

$$P_{av} = \int_0^b \int_0^a \frac{\tilde{\beta} C_1^2}{2\omega\mu} \sin^2\left(\frac{m\pi}{a}x\right) dx dy = \frac{\tilde{\beta} C_1^2 ab}{4\omega\mu} \quad [p.277]$$

 For the parallel plate TEM (TM<sub>0</sub>) mode the total power through the guide is

$$P_{av} = \frac{1}{2} \eta C_4^2 ba \quad [p.277]$$

Dielectric slab waveguide

TM Modes The non-zero field components are  $E_z$ ,  $E_x$ , and  $H_y$

For  $x \leq -d/2$

$$E_z^0(x) = C_o \sin(\beta_x x) + C_e \cos(\beta_x x) \quad [4.34]$$

where the transverse propagation constant is given by

$$\beta_x^2 = \omega^2 \mu_d \epsilon_d - \bar{\beta}^2 = h_d^2 \quad [4.35]$$

For  $x \geq d/2$

$$E_z^0(x) = \begin{cases} C_a e^{-\alpha_x(x-d/2)}, & x \geq d/2 \\ C_a e^{\alpha_x(x+d/2)}, & x \leq -d/2 \end{cases} \quad [4.36]$$

where the transverse attenuation constant is given by

$$\alpha_x^2 = \bar{\beta}^2 - \omega^2 \mu_0 \epsilon_0 = -h_0^2 \text{ or } \alpha_x = \sqrt{\bar{\beta}^2 - \omega^2 \mu_0 \epsilon_0} \quad [4.37]$$

The cutoff frequencies are given by

$$f_{c_{TMm}} = \frac{(m-1)}{2d\sqrt{\mu_d \epsilon_d - \mu_0 \epsilon_0}} \quad \begin{matrix} m = 1, 3, 5, \dots \text{odd} \\ m = 2, 4, 6, \dots \text{even} \end{matrix} \quad [4.45]$$

TE Modes The non-zero field components are  $H_z$ ,  $H_x$ , and  $E_y$

For  $x \leq -d/2$

$$H_z^0(x) = C_o \sin(\beta_x x) + C_e \cos(\beta_x x) \quad [4.46]$$

where the transverse propagation constant

$$\beta_x^2 = \omega^2 \mu_d \epsilon_d - \bar{\beta}^2 = h_d^2$$

For  $x \geq d/2$

$$H_z^0(x) = \begin{cases} C_a e^{-\alpha_x(x-d/2)}, & x \geq d/2 \\ C_a e^{\alpha_x(x+d/2)}, & x \leq -d/2 \end{cases}$$

where the transverse propagation constant

$$\alpha_x^2 = \bar{\beta}^2 - \omega^2 \mu_0 \epsilon_0 = -h_0^2 \text{ or } \alpha_x = \sqrt{\bar{\beta}^2 - \omega^2 \mu_0 \epsilon_0} \quad [4.37]$$

For Odd TM Modes:

$$\frac{\alpha_x}{\beta_x} = \frac{\epsilon_0}{\epsilon_d} \tan\left(\frac{\beta_x d}{2}\right) \quad [4.40]$$

For Even TM Modes:

$$\frac{\alpha_x}{\beta_x} = -\frac{\epsilon_0}{\epsilon_d} \cot\left(\frac{\beta_x d}{2}\right) \quad [4.44]$$

For ALL TM Modes:

$$\alpha_x = \sqrt{\omega^2 (\mu_d \epsilon_d - \mu_0 \epsilon_0) - \beta_x^2} \quad [4.42]$$

The cutoff frequencies are given by

$$f_{c_{TMm}} = \frac{(m-1)}{2d\sqrt{\mu_d \epsilon_d - \mu_0 \epsilon_0}} \quad \begin{matrix} m = 1, 3, 5, \dots \text{odd} \\ m = 2, 4, 6, \dots \text{even} \end{matrix} \quad [4.49]$$

## Dielectric covered ground plane

$$f_{c_{TMm}} = \frac{(m-1)}{2d\sqrt{\mu_d \epsilon_d - \mu_0 \epsilon_0}} \quad \begin{matrix} m = 1, 3, 5, \dots \text{odd} \text{ } _{TM_m} \\ m = 2, 4, 6, \dots \text{even} \text{ } _{TE_m} \end{matrix} \quad [4.50]$$

## Dielectric slab waveguide ray theory

$$\tan \theta_i = \frac{\bar{\beta}}{\beta_x} \quad [\text{p.306}]$$

EECS 412 Electromagnetic Fields III Formula Sheet

Detailed example of odd TM Modes for slab dielectric waveguide: free space above the guide

For  $x \geq d/2$

$$E_z^0(x) = C_0 \sin\left(\frac{\beta_x d}{2}\right) e^{-\alpha_x(x-d/2)} \quad [4.39a]$$

$$E_x^0(x) = -\frac{j\bar{\beta}}{\alpha_x} C_0 \sin\left(\frac{\beta_x d}{2}\right) e^{-\alpha_x(x-d/2)} \quad [4.39b]$$

$$H_y^0(x) = -\frac{j\omega\epsilon_0}{\alpha_x} C_0 \sin\left(\frac{\beta_x d}{2}\right) e^{-\alpha_x(x-d/2)} \quad [4.39c]$$

For  $|x| \leq d/2$

$$E_z^0(x) = C_0 \sin(\beta_x x) \quad [4.39d]$$

$$E_x^0(x) = -\frac{j\bar{\beta}}{\beta_x} C_0 \cos(\beta_x x) \quad [4.39e]$$

$$H_y^0(x) = \frac{j\omega\epsilon_0}{\beta_x} C_0 \cos(\beta_x x) \quad [4.39f]$$

For  $x \leq -d/2$

$$E_z^0(x) = C_0 \sin\left(\frac{\beta_x d}{2}\right) e^{-\alpha_x(x+d/2)} \quad [4.39g]$$

$$E_x^0(x) = -\frac{j\bar{\beta}}{\alpha_x} C_0 \sin\left(\frac{\beta_x d}{2}\right) e^{-\alpha_x(x+d/2)} \quad [4.39h]$$

$$H_y^0(x) = -\frac{j\omega\epsilon_0}{\alpha_x} C_0 \sin\left(\frac{\beta_x d}{2}\right) e^{-\alpha_x(x+d/2)} \quad [4.39i]$$